

On the Cauchy Problem for a Nonlinearly Dispersive Wave Equation

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Abstract

We establish the local well-posedness for a new nonlinearly dispersive wave equation and we show that the equation has solutions that exist for indefinite times as well as solutions which blowup in finite time. Furthermore, we derive an explosion criterion for the equation and we give a sharp estimate from below for the existence time of solutions with smooth initial data.

1 Introduction

In this letter we consider the Cauchy problem for the nonlinear equation:

$$\begin{aligned} u_t - u_{txx} + 2\omega u_x + 3uu_x &= \gamma(2u_x u_{xx} + uu_{xxx}), & t > 0, \quad x \in \mathbb{R}, \\ u(0, x) &= u_0(x), & x \in \mathbb{R}, \end{aligned} \quad (1.1)$$

with $\omega \geq 0$ and $\gamma \in \mathbb{R}$ fixed constants.

With $\gamma = 1$ in equation (1.1) we find the Camassa–Holm equation for the unidirectional propagation of shallow water waves over a flat bottom, $u(t, x)$ standing for the fluid velocity at time $t \geq 0$ in the spatial x direction and ω being a nonnegative parameter related to the critical shallow water speed (see [3] and [16]). In this case equation (1.1) has a bi-Hamiltonian structure [14] and is completely integrable (see [1, 6, 18]). The solitary waves are smooth if $\omega > 0$, they become peaked in the limiting case $\omega = 0$, and interact like solitons [3].

For $\gamma = 0$, equation (1.1) becomes the regularized long wave equation, a well-known equation for surface waves in a channel [2]. All solutions are global and the solitary waves are smooth. Despite having a Hamiltonian structure, the equation is not integrable and its solitary waves are not solitons [13].

For $\omega = 0$ and $\gamma \in \mathbb{R}$, equation (1.1) particularizes to a nonlinear dispersive model for finite-length and small-amplitude radial deformation waves in thin cylindrical compressible hyperelastic rods (see [10] and [11]), $u(t, x)$ representing the radial stretch relative to a pre-stressed state. The case $\gamma = 1$ (limiting case of the Camassa–Holm equation) has been extensively studied: the equation has global solutions [5, 7] and also solutions which blow-up in finite time [5, 8, 7]. It is an integrable Hamiltonian system [1, 6] and its solitary waves are peaked and interact like solitons [3].

For $\omega > 0$ and $\gamma \in \mathbb{R}$, equation (1.1) seems not yet to have been discussed. The aim of this letter is to prove local well-posedness of strong solutions to equation (1.1) for a large class of initial data, to analyze global existence of smooth solutions, as well as to get a blow-up criterion and a sharp estimate from below of the existence time for strong solutions.

2 Local Well-posedness

To prove the local existence of solutions for equation (1.1) we will apply Kato's semi-group approach [17]. For this purpose, let us reformulate equation (1.1) as a quasi-linear evolution equation for $y := u - u_{xx}$,

$$y_t + \gamma y_x u + 2\gamma y u_x + 2\omega u_x + 3(1 - \gamma)u u_x = 0, \quad t > 0, \quad x \in \mathbb{R}. \quad (2.1)$$

Note that if $p(x) := \frac{1}{2}e^{-|x|}$, $x \in \mathbb{R}$, then $(1 - \partial_x^2)f = p * f$ for all $f \in L^2(\mathbb{R})$ so $p * y = u$, where $*$ denotes convolution. Using this identity, we can rewrite equation (2.1) as follows,

$$u_t + \gamma u u_x + \partial_x p * \left(\frac{3 - \gamma}{2} u^2 + \frac{\gamma}{2} u_x^2 + 2\omega u \right) = 0, \quad t > 0, \quad x \in \mathbb{R}. \quad (2.2)$$

Equation (2.2) is suitable for applying Kato's theory. A framework that is similar to those provided in [7] and [20] can be constructed. As an outcome, we obtain

Theorem 1. *Given $u_0 \in H^s(\mathbb{R})$, $s > \frac{3}{2}$, there exists a maximal $T = T(\gamma, \omega, u_0) > 0$, and a unique solution u to equation (1.1), such that*

$$u = u(\cdot, u_0) \in C([0, T]; H^s(\mathbb{R})) \cap C^1([0, T]; H^{s-1}(\mathbb{R})).$$

Moreover, the solution depends continuously on the initial data.

Theorem 1 contains as particular cases (for various choices of the constants γ and ω) the well-posedness results obtained in [7, 15, 19] and [20].

3 Global solutions

To answer affirmatively the question whether there are solutions of equation (1.1) that exist indefinitely in time, we prove that for every choice of $\gamma \in \mathbb{R}$ and $\omega > 0$ there are smooth solitary wave solutions $u(t, x) = \varphi_c(x - ct)$ propagating at certain speeds $c > 0$.

If $u(t, x) = \varphi_c(x - ct)$ is a smooth solitary wave solution of equation (1.1), then

$$-c\varphi'_c + c\varphi_c''' + 3\varphi'_c\varphi_c + 2\omega\varphi'_c = 2\gamma\varphi'_c\varphi_c'' + \gamma\varphi_c\varphi_c'''. \quad (3.1)$$

Fix $x \in \mathbb{R}$ and integrate equation (3.1) on $(-\infty, x]$ to obtain

$$(2\omega - c)\varphi_c + c\varphi_c'' + \frac{3}{2}(\varphi_c)^2 = \frac{\gamma}{2}(\varphi_c')^2 + \gamma\varphi_c\varphi_c''. \quad (3.2)$$

Multiplying equation (3.2) by φ_c' , an integration on $(-\infty, x]$ leads to

$$(\varphi_c')^2(c - \gamma\varphi_c) = \varphi_c^2(c - 2\omega - \varphi_c).$$

A detailed analysis of the above ordinary differential equation yields the following result.

Theorem 2. *Let $\omega > 0$ and $\gamma \in \mathbb{R}$. Then for every $c > 0$ satisfying $c(\gamma - 1) < 2\omega\gamma$ and $c > 2\omega$, there exists a nontrivial smooth solitary wave solution $u(t, x) = \varphi_c(x - ct)$ of equation (1.1) and its profile φ_c decays exponentially to zero at infinity.*

4 Solutions with finite life-span

In this section we address the question of the formation of singularities for solutions to equation (1.1). It turns out that there are smooth initial data for which the corresponding solution to equation (1.1) does not exist globally in time. The precise blowup scenario is described. We also give an explosion criterion and a sharp estimate from below for the maximal existence time of smooth solutions to equation (1.1). Analogous to the case of the Camassa–Holm equation [4], equation (1.1) can be written in Hamiltonian form and has the invariant

$$E(u) = \int_{\mathbb{R}} (u^2 + u_x^2) dx. \quad (4.1)$$

The invariance of $E(u)$ ensures that all solutions of equation (1.1) are uniformly bounded as long as they exist. Arguments similar to those developed in [9] for the elasticity model can be pursued to obtain the following result.

Theorem 3. *Given $u_0 \in H^s(\mathbb{R})$, $s > \frac{3}{2}$, the solution $u = u(\cdot, u_0)$ of equation (1.1) is uniformly bounded. Blow up in finite time $T < +\infty$ occurs if and only if*

$$\liminf_{t \uparrow T} \left\{ \inf_{x \in \mathbb{R}} [\gamma u_x(t, x)] \right\} = -\infty.$$

As an interesting outcome of the previous result, note that in the special case $\gamma = 0$, all solutions to equation (1.1) are global. Also, in the particular cases when equation (1.1) models water waves or waves on an elastic rod, the previously described blowup corresponds physically to a breaking wave, respectively rod.

We discuss now the question of finite time blow-up of solution to equation (1.1) with rather general initial data. Let $T > 0$ be the maximal existence time of the solution $u(t, \cdot)$ of equation (1.1) with initial data $u_0 \in H^s(\mathbb{R})$, $s > \frac{3}{2}$. Differentiate equation (2.2) with respect to x , and multiply the obtained equation by γ . Since $\partial_x^2 p * f = p * f - f$, we have

$$\begin{aligned} \gamma u_{tx} = & -\frac{\gamma^2}{2}u_x^2 - \gamma^2 u u_{xx} + \frac{(3-\gamma)\gamma}{2}u^2 \\ & + 2\omega\gamma u - p * \left(\frac{(3-\gamma)\gamma}{2}u^2 + \frac{\gamma^2}{2}u_x^2 + 2\omega\gamma u \right). \end{aligned} \quad (4.2)$$

Define $m(t) = \gamma u_x(t, \xi(t)) = \inf_{x \in \mathbb{R}} \{\gamma u_x(t, x)\}$, where $\xi(t)$ is any point where this minimum is attained. Since we deal with a minimum, $u_{xx}(t, \xi(t)) = 0$ for all $t \in [0, T)$. On the other hand, the regularity of the solution $u(t, x)$ provided by Theorem 1 allows us to infer from the main result in [8] that the function $m(t)$ is almost everywhere differentiable on $(0, T)$, with

$$\frac{dm}{dt} = \gamma u_{tx}(t, \xi(t)) \quad \text{a.e. on } (0, T).$$

Hence we obtain that

$$\begin{aligned} m'(t) = & -\frac{1}{2}m^2(t) + \frac{(3-\gamma)\gamma}{2}u^2(t, \xi(t)) + 2\omega\gamma u(t, \xi(t)) \\ & - p * \left(\frac{(3-\gamma)\gamma}{2}u^2 + \frac{\gamma^2}{2}u_x^2 + 2\omega\gamma u \right) (t, \xi(t)). \end{aligned} \quad (4.3)$$

Involved manipulations of the equations (4.2) and (4.3) and estimates analogous to those in [12] for the special case of the Camassa–Holm equation lead us to the following result.

Theorem 4. (i) Given $u_0 \in H^s(\mathbb{R})$, $s > \frac{3}{2}$, and assume that we can find $x_0 \in \mathbb{R}$ with

$$\gamma u'_0(x_0) < - \left(\frac{|\gamma - 3|\gamma|}{2} E(u_0) + 4\sqrt{2}\omega|\gamma|(E(u_0))^{1/2} \right)^{1/2}, \quad \gamma \neq 0.$$

Then the corresponding solution to equation (1.1) blows up in finite time.

(ii) Given $u_0 \in H^s(\mathbb{R})$, $s > \frac{3}{2}$, the maximal existence time of the solution to equation (1.1) with initial data u_0 is estimated from below by

$$T_{u_0} = \begin{cases} \frac{-2 \arctan \frac{\left(\frac{(3-\gamma)\gamma}{2} E(u_0) + 4\sqrt{2}\omega|\gamma|(E(u_0))^{1/2}\right)^{1/2}}{m(0)}}{\left(\frac{(3-\gamma)\gamma}{2} E(u_0) + 4\sqrt{2}\omega|\gamma|(E(u_0))^{1/2}\right)^{1/2}}, & \text{if } 0 < \gamma < \frac{3}{2}, \\ \frac{-2 \arctan \frac{\left(\frac{\gamma^2}{2} E(u_0) + 4\sqrt{2}\omega|\gamma|(E(u_0))^{\frac{1}{2}}\right)^{1/2}}{m(0)}}{\left(\frac{\gamma^2}{2} E(u_0) + 4\sqrt{2}\omega|\gamma|(E(u_0))^{1/2}\right)^{1/2}}, & \text{if } \frac{3}{2} \leq \gamma \leq 3, \\ \frac{-2 \arctan \frac{\left(\frac{(2\gamma-3)\gamma}{2} E(u_0) + 4\sqrt{2}\omega|\gamma|(E(u_0))^{1/2}\right)^{1/2}}{m(0)}}{\left(\frac{(2\gamma-3)\gamma}{2} E(u_0) + 4\sqrt{2}\omega|\gamma|(E(u_0))^{1/2}\right)^{1/2}}, & \text{if } \gamma > 3 \text{ or } \gamma < 0, \end{cases}$$

where $m(0) := \inf_{x \in \mathbb{R}} \{\gamma \partial_x u_0(x)\}$. The previous estimate is sharp.

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